

# Experimental Determination of the Heat of Dissociation of $\text{N}_2\text{O}_4 \rightarrow 2\text{NO}_2$ from the Temperature Dependence of Absolute Infrared Intensities\*

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The heat of dissociation of  $\text{N}_2\text{O}_4$  has been determined from measurements of the temperature dependence of the integrated absorption of  $\text{NO}_2$  and  $\text{N}_2\text{O}_4$  vibration-rotation bands. Results have been obtained that are in acceptable accord with the earlier estimate of Giauque and Kemp.

ABSOLUTE infrared intensity measurements have been performed by using the technique of self-broadening<sup>1</sup> on a number of vibration-rotation bands belonging to  $\text{NO}_2$  and  $\text{N}_2\text{O}_4$ . A complete account of the experimental results is given elsewhere<sup>2</sup>; the experimental data were obtained on the equilibrium gases formed over liquid  $\text{N}_2\text{O}_4$ .

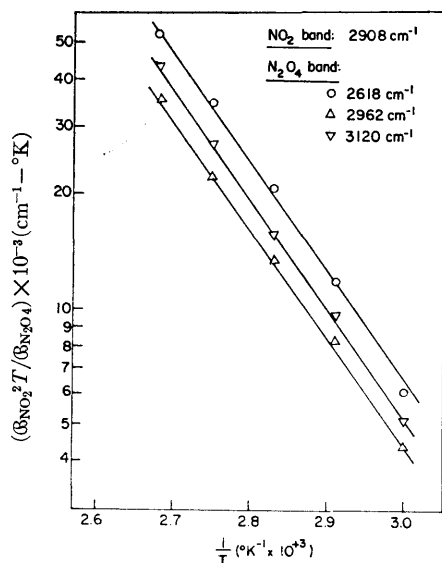


FIG. 1. Plots for determining  $\Delta H^\circ$  for the reaction  $\text{N}_2\text{O}_4 \rightarrow 2\text{NO}_2$  from experimental values of the parameters  $B$ . The three curves were obtained using the strongest  $\text{N}_2\text{O}_4$  bands in the region from 1 to 5 microns ( $\text{NO}_2$  band:  $2908 \text{ cm}^{-1}$ ;  $\text{N}_2\text{O}_4$  bands:  $\circ$   $2618 \text{ cm}^{-1}$ ;  $\triangle$   $2962 \text{ cm}^{-1}$ ;  $\nabla$   $3120 \text{ cm}^{-1}$ ).

The measured temperature dependence of the integrated intensities may be used to determine the heat of dissociation. Using standard procedures, we define the quantity

$$B = \int_{\omega_1}^{\omega_2} \left[ \ln \frac{I_{\omega}^{o,a}}{I_{\omega}^a} \right] d\omega, \quad (1)$$

where  $I_{\omega}^{o,a}$  and  $I_{\omega}^a$  denote, respectively, the incident and transmitted apparent flux densities at the wave

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<sup>1</sup> S. S. Penner and D. Weber, *J. Chem. Phys.* **19**, 817 (1951).

<sup>2</sup> A. Guttman, Ph.D. thesis, California Institute of Technology, Pasadena, 1962.

number  $\omega$ . The integration limits in Eq. (1) are chosen in such a way that an entire vibration-rotation band is included between  $\omega_1$  and  $\omega_2$ . It is well known that, if the spectral absorption coefficient is constant over the wave number range corresponding to the effective slit width, then

$$dB/dX = \alpha, \quad (2)$$

where  $X = pl$  is the optical depth in cm-atm if  $p$  is the partial pressure of absorber in atm and  $l$  denotes the geometric length in cm;  $\alpha$  represents the integrated intensity of the vibration-rotation band in  $\text{cm}^{-2}\text{-atm}^{-1}$  at the temperature at which the experiment was performed.

If we consider two bands, one belonging to  $\text{NO}_2$  and another belonging to  $\text{N}_2\text{O}_4$ , then it is apparent from Eq. (2) that

$$p_{\text{NO}_2} = B_{\text{NO}_2} / \alpha_{\text{NO}_2} l$$

and

$$p_{\text{N}_2\text{O}_4} = B_{\text{N}_2\text{O}_4} / \alpha_{\text{N}_2\text{O}_4} l,$$

provided adequate broadening of the spectral lines has been achieved. For the small temperature range covered in the experimental studies and extending from about  $295^\circ\text{K}$  to about  $395^\circ\text{K}$ ,<sup>3</sup> we may assume, in good approximation, that the integrated intensities are inversely proportional to the absolute temperatures  $T$ . Hence

$$\alpha_{\text{NO}_2} = \alpha_{o,\text{NO}_2} (T_o/T)$$

and

$$\alpha_{\text{N}_2\text{O}_4} = \alpha_{o,\text{N}_2\text{O}_4} (T_o/T),$$

where the subscript  $o$  identifies suitable reference conditions. The preceding expressions lead to the following relation for the equilibrium constant at constant pressure:

$$K_p = \frac{p_{\text{NO}_2}^2}{p_{\text{N}_2\text{O}_4}} = \left( \frac{\alpha_{o,\text{N}_2\text{O}_4}}{\alpha_{o,\text{NO}_2}^2 T_o l} \right) \left( \frac{B_{\text{NO}_2}^2 T}{B_{\text{N}_2\text{O}_4}} \right). \quad (3)$$

For specified reference conditions and constant geometric length  $l$ , the first factor in Eq. (3) is independent

<sup>3</sup> The corresponding equilibrium vapor pressures of  $\text{NO}_2$  over liquid  $\text{N}_2\text{O}_4$  lie between about 1.5 and 25.0 atm; see H. H. Reamer and B. H. Sage, *Ind. Eng. Chem.* **44**, 185 (1952).

of temperature. Hence it follows from Eq. (3) that

$$\frac{d \ln K_p}{d(1/T)} = -\frac{\Delta H^\circ}{R} = \frac{d \ln (\mathfrak{B}_{\text{NO}_2^2 T} / \mathfrak{B}_{\text{N}_2\text{O}_4})}{d(1/T)}, \quad (4)$$

where  $\Delta H^\circ$  is the molar heat of dissociation of  $\text{N}_2\text{O}_4$  and  $R$  represents the molar gas constant.

Experimentally determined values of the quantity  $\mathfrak{B}_{\text{NO}_2^2 T} / \mathfrak{B}_{\text{N}_2\text{O}_4}$  for the  $\text{NO}_2$  band centered at  $2908 \text{ cm}^{-1}$  and for three  $\text{N}_2\text{O}_4$  bands (centered at 2618, 2962, and  $3120 \text{ cm}^{-1}$ ) are plotted in Fig. 1 as a function of  $1/T$ . The geometric length  $l$  was set at  $0.124 \text{ cm}$  by using a suitable spacer.<sup>2</sup> The data shown in Fig. 1 yield values of  $\Delta H^\circ = 13\,070$ ,  $12\,980$ , and  $13\,160 \text{ cal/mole}$  of  $\text{N}_2\text{O}_4$  for the curves corresponding to the 2618, 2962, and  $3120 \text{ cm}^{-1}$  bands of  $\text{N}_2\text{O}_4$ , respectively. These results are in reasonably satisfactory agreement with the value

$\Delta H^\circ = 13\,693 \pm 14 \text{ cal/(mole of } \text{N}_2\text{O}_4)$  reported by Giaque and Kemp<sup>4</sup> at  $298^\circ\text{K}$ .

The preceding results provide an important check on the self-consistency of the absolute intensity measurements although they have no direct bearing on the absolute accuracy of the results. The procedure for determining heats of reaction by measuring the temperature dependence of (absolute) infrared intensities is clearly applicable also to other molecules. An experiment similar to that described in the present note was suggested to one of us (S.S.P.) by Professor Farrington Daniels in 1942.

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<sup>4</sup> W. F. Giaque and J. D. Kemp, *J. Chem. Phys.* **6**, 40 (1938). It is easily shown that the equation of state proposed by Giaque and Kemp leads to cancellation of real gas effects in the expression for the equilibrium constant.